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# A high performance magnetic fluid based on carbon modified magnetite (Fe $_3O_4$ ) nanospheres

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#### ABSTRACT

This work reported a carbon layer modified magnetite (Fe<sub>3</sub>O<sub>4</sub>) nanospheres (Fe<sub>3</sub>O<sub>4</sub>@C) which could be applied in preparing high performance magnetic fluid. The porous Fe<sub>3</sub>O<sub>4</sub>@C nanospheres were prepared by calcinating the Fe<sub>3</sub>O<sub>4</sub>@Polydopamine (Fe<sub>3</sub>O<sub>4</sub>@PDA) precursor in N<sub>2</sub> atmosphere. The shear stress of Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid was as high as 24.9 Pa at saturated magnetic field, which was about 3.8 times larger than pristine Fe<sub>3</sub>O<sub>4</sub>-based magnetic fluid (6.6 Pa). Furthermore, the particle level dynamic simulation was employed to analyze the microstructure evolution and shear stress of Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid under applying an external magnetic field. The chain-like model well explained the magnetorheological (MR) mechanism and the simulation results matched perfectly with the experimental data. Since the crystallinity of Fe<sub>3</sub>O<sub>4</sub> was improved during the pyrolysis, Fe<sub>3</sub>O<sub>4</sub>@C nanospheres possessed a higher saturation magnetization. The stronger magnetic fluid was resulted in a much steadier Fe<sub>3</sub>O<sub>4</sub>@C chain-like structure. As a result, it was found that the superior MR performance and high stability of the Fe<sub>3</sub>O<sub>4</sub>@C chased magnetic fluid was reported a core-shell and porous structures of Fe<sub>3</sub>O<sub>4</sub>@C. This work proposed a valuable magnetic particle model to prepare high performance magnetic fluid and further understand the origination of MR effects.

#### 1. Introduction

The magnetic fluids are a category of intelligent materials that compose of magnetic particles and liquid dispersed medium, and their viscosities can be tuned quickly and reversibly by applying the external magnetic field [1]. Because of the magnetic dipole-dipole interactions, the magnetic particles form a chain-like microstructure along magnetic field direction, which lead the magnetic fluids to be transformed from fluid-like to solid-like state. So, the mechanical properties of magnetic fluids change significantly under the external magnetic field [2,3]. Due to the superior rheological behaviour and controllable mechanical properties, magnetic fluids have received a great deal of focuses in diverse application fields, such as intelligent sensors [4], dampers [5,6], exercise industry [7,8] and biological field [9,10], etc.

During the past decades, massive efforts have been made to develop the high-performance magnetic fluids. A variety of magnetic particles, such as ferrite [11–13], carbonyl iron [14,15], chromium dioxide [16] and  $Fe_3O_4$  [17,18] have been used in magnetic fluids. It is generally

cognized that MR effect results from the transition of magnetic particles-assembly-structure from disordered to ordered under the external magnetic field. So the morphology, magnetic property, size and innerstructure of magnetic particles exhibit critical influence on the rheological properties [19]. Magnetic particles with various shape such as cubic [20], octahedral [21], rodlike [22] and spherical [23,24] have been developed for magnetic fluids. In comparison to the spherical particles, the anisotropic particles can assemble into stronger chains structure because of their larger friction force [25]. Furthermore, recent studies have manifested that the magnetically induced bundle wire-like structures originated from the magnetic chains were more stable than that of the magnetic nanospheres, therefore the magnetic chains based magnetic fluid possessed preferable MR effect [26]. The inner-structure of magnetic particles plays a one of critical roles for the MR properties. Owing to the low intensity, the hollow magnetic particles endow the final magnetic fluid with high stability. The friction, blocking, and water encapsulation in the hollow structure also strengthen the chainlike structure, thus the hollow magnetic particles show higher MR

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performance [27,28]. Moreover, the particles size of magnetic particles also affect the mechanical properties of the magnetic fluid. The research indicates that MR effects enhanced with increasing of the particle size and surface interaction [29]. To this end, it is meaningful to investigate the effect of particle nanostructure on the rheological properties.

Recently, the Fe<sub>3</sub>O<sub>4</sub>-based magnetic fluids have received increasing interests due to its compatible surface, adjustable size and low density which match well with the carrying fluid [30]. The Fe<sub>3</sub>O<sub>4</sub> with different size, shape, magnetic property and surface state can be prepared conveniently [29,31–33]. Thus the Fe<sub>3</sub>O<sub>4</sub> based magnetic fluid is beneficial to understanding the mechanism of the MR effects by easily tuning the effects of various parameters of magnetic particles [34]. Hence, the investigation of Fe<sub>3</sub>O<sub>4</sub> based magnetic fluid is greatly attractive not only for its widely application but also as an ideal candidate for comprehending the MR mechanism. Unfortunately, in previous study, few works have been reported that how the surface state of Fe<sub>3</sub>O<sub>4</sub> particles influence the mechanism of magnetic fluid and prepare high performance magnetic fluid, it is essential to study the magnetic fluid based on Fe<sub>3</sub>O<sub>4</sub> particles with different inner-structure.

In this work, the Fe<sub>3</sub>O<sub>4</sub>@C nanospheres, which were synthesized *via* pyrolyzing the Fe<sub>3</sub>O<sub>4</sub>@PDA core-shell nanospheres, could be dispersed into the aqueous solution to prepare high performance magnetic fluid. The rheological testing indicated that the Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid possessed high MR effect and good stability. In comparison to pristine Fe<sub>3</sub>O<sub>4</sub> and Fe<sub>3</sub>O<sub>4</sub>@PDA particles, the Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid presented a higher MR performance. Furthermore, the particle level dynamic simulation was applied to analyse the magnetic responds of the magnetic fluid based on Fe<sub>3</sub>O<sub>4</sub>@C nanospheres. The simulated chain-like model and shear stress as a function of magnetic field agreed well with the experimental data. Finally, a feasible mechanism was proposed to analyse the reasons for the enhancement MR performance of Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid. This work supplied valuable information for designing high performance magnetic fluid and understanding the MR mechanism.

#### 2. Results and discussion

#### 2.1. Characterization of magnetic Fe<sub>3</sub>O<sub>4</sub>@C nanospheres

The synthesis procedure sketch of Fe<sub>3</sub>O<sub>4</sub>@C is exposited in Fig. 1a. Firstly, the Fe<sub>3</sub>O<sub>4</sub> nanospheres are prepared via a solvothermal reaction [35]. The transmission electronic microscopy (TEM) and scanning electron microscope (SEM) images in Fig. 1b and c indicate that Fe<sub>3</sub>O<sub>4</sub> nanoparticles are monodispersed with uniform spherical morphology. The average size of Fe<sub>3</sub>O<sub>4</sub> is about 100 nm. From SEM image of Fe<sub>3</sub>O<sub>4</sub>, it is obviously seen that the surface of Fe<sub>3</sub>O<sub>4</sub> is very rough because the Fe<sub>3</sub>O<sub>4</sub> nanosphere is composed of numerous small nanocrystals and exhibits a typical cluster-like nanostructure. Then, the PDA shell are equably coated on the surface of Fe<sub>3</sub>O<sub>4</sub> through the simple oxidation polymerization, due to its brilliant ability in adhesion to both organic and inorganic substrates. After the decoration of PDA layer, the Fe<sub>3</sub>O<sub>4</sub>@ PDA composite materials still maintain a monodisperse spherical morphology (Fig. 1d and e). From the high magnification TEM image of  $Fe_3O_4$ @PDA (inset of Fig. 1d), there is an evident pale layer located at the periphery of Fe<sub>3</sub>O<sub>4</sub>. Additionally, the surface of Fe<sub>3</sub>O<sub>4</sub>@PDA (Fig. 1e) is smoother than that of  $Fe_3O_4$  and the average size of  $Fe_3O_4@$ PDA nanospheres increases to about 110 nm. All of results prove the existence of uniform PDA layer. After PDA shell being carbonized to carbon shell, the Fe<sub>3</sub>O<sub>4</sub>@PDA transform to Fe<sub>3</sub>O<sub>4</sub>@C nanospheres. The Fe<sub>3</sub>O<sub>4</sub>@C nanospheres show spherical morphology without apparent aggregation, which gives the credit to the presence of PDA shell (Fig. 1f and g). Interestingly, the average size of Fe<sub>3</sub>O<sub>4</sub>@C is about 100 nm, which is a little smaller than the Fe<sub>3</sub>O<sub>4</sub>@PDA precursor. Form the high magnification TEM image (inset of Fig. 1f), it can be observed that the Fe<sub>3</sub>O<sub>4</sub>@C nanospheres possess porous structures. The BrunauerEmmett-Teller (BET) surface area of Fe<sub>3</sub>O<sub>4</sub>@C is 66.8 m<sup>2</sup>/g, which is higher than that of Fe<sub>3</sub>O<sub>4</sub> (50.8 m<sup>2</sup>/g) (Fig. S2). After carefully comparing the TEM images of Fe<sub>3</sub>O<sub>4</sub> and Fe<sub>3</sub>O<sub>4</sub>@C (Fig. 1b and f), it can be observed that the Fe<sub>3</sub>O<sub>4</sub> nanograins sizes get bigger after the pyrolysis process.

X-ray photoelectron spectroscopy (XPS) is performed to examine the shell structure of the final products. As exhibited in Fig. 2a, the signals of Fe, O and C elements are clearly detected in the XPS spectrum of Fe<sub>3</sub>O<sub>4</sub>, which are originated from the Fe<sub>3</sub>O<sub>4</sub> and residual polymer reagents. After being covered by PDA shell, the specific signal of N 1s is found and the intensity of C 1s is enhanced in the corresponding curve, demonstrating that PDA layer is wrapped on the Fe<sub>3</sub>O<sub>4</sub> core successfully. Furthermore, Fe 2p and Fe 3p signals still exist but the signal intensities of them are reduced significantly, thus the thickness of PDA layer is less than 10 nm because the detection thickness of the XPS is about 10 nm, which is in accordance with the TEM analysis. The C1s core-level spectrum of Fe<sub>3</sub>O<sub>4</sub>@PDA shows four characteristic peaks at 288.75, 287.9, 286.15 and 284.8 eV, which can be assigned to O-C=O, C = O, C-N, C = C, respectively (Fig. 2b). The N 1s spectrum can be split to three peaks at 401.95, 400.1 and 398.45 eV (Fig. 2c). The first peak at 401.95 eV corresponds to the amine group (R-NH<sub>2</sub>), which must come from the small amount of physically self-assembled dopamine. The peaks at 400.1 and 398.45 eV can be ascribed to the substituted amine (-N-H) and imino (=N-R) groups of PDA. The wide-scan XPS spectrum of Fe<sub>3</sub>O<sub>4</sub>@C is similar to that of Fe<sub>3</sub>O<sub>4</sub>@PDA, because the Fe<sub>3</sub>O<sub>4</sub>@C is obtained by the pyrolysis of Fe<sub>3</sub>O<sub>4</sub>@PDA. Nevertheless, the C 1s and N 1s core-level spectra of Fe<sub>3</sub>O<sub>4</sub>@C (Fig. 2b and c) is slightly different from that of Fe<sub>3</sub>O<sub>4</sub>@PDA. After pyrolysis, the intensity of C = Cand imino (=N-R) peaks are apparently increased, which can be attributed to the formation of graphite carbon.

The Raman spectra further prove the presence of PDA and C layer. As shown in Fig. 3a, there are two visible peaks at 1582 and 1368 cm<sup>-1</sup>, corresponding to the G and D bands, respectively. After the pyrolysis, the crystallization of the carbon component in the Fe<sub>3</sub>O<sub>4</sub>@C is better than the Fe<sub>3</sub>O<sub>4</sub>@PDA, thus sharper peaks are found in the Raman spectra. Moreover, it can be distinctly observed that the weight loss of Fe<sub>3</sub>O<sub>4</sub>@C (9.9%) is much lower than that of Fe<sub>3</sub>O<sub>4</sub>@PDA (20.6%) and higher than that of Fe<sub>3</sub>O<sub>4</sub> (6.8%) (Fig. 3b). All results indicate that there is a thin carbon layer on the surface of Fe<sub>3</sub>O<sub>4</sub> in the Fe<sub>3</sub>O<sub>4</sub>@C.

The XRD patterns are applied to track the synthetic process of the Fe<sub>3</sub>O<sub>4</sub>@C nanospheres. As shown in Fig. 3c, all the products show the pure face-cubic centred magnetite phase of the Fe<sub>3</sub>O<sub>4</sub>, because the crystal phase in Fe<sub>3</sub>O<sub>4</sub>, Fe<sub>3</sub>O<sub>4</sub>@PDA and Fe<sub>3</sub>O<sub>4</sub>@C is Fe<sub>3</sub>O<sub>4</sub>. Based on the Scherrer formula, the nanograins sizes of Fe<sub>3</sub>O<sub>4</sub> and Fe<sub>3</sub>O<sub>4</sub>@C are calculated to be 10 and 15 nm, respectively. This result indicates the crystallinity of Fe<sub>3</sub>O<sub>4</sub> is critically improved during the pyrolysis process, which is keeping with the TEM images analyses.

The magnetic properties of as-prepared samples are evaluated by SQUID at room temperature. The hysteresis loop of each sample presents the typical superparamagnetic characteristic, of which the coercivity (Hc) and remanent magnetization are almost zero (Fig. 3d). Therefore, these nanopheres are easy to be dispersed to form homogenous solution. All samples can be completely saturated at an external magnetic field of 20 KOe. The saturation magnetizations of Fe<sub>3</sub>O<sub>4</sub>, Fe<sub>3</sub>O<sub>4</sub>@PDA and Fe<sub>3</sub>O<sub>4</sub>@C are 69.8, 52.6 and 69.8 emu/g, respectively. There is no doubt that the magnetic properties of each sample are derived from the Fe<sub>3</sub>O<sub>4</sub> core, because the PDA and C shell are non-magnetic. The introduction of non-magnetic material reduces the saturation magnetization of the target product but the saturation magnetization of Fe<sub>3</sub>O<sub>4</sub>@C is as high as original Fe<sub>3</sub>O<sub>4</sub>, indicating that the saturation magnetization of Fe<sub>3</sub>O<sub>4</sub> component in Fe<sub>3</sub>O<sub>4</sub>@C is higher than original Fe<sub>3</sub>O<sub>4</sub>. Combining the TG analysis, the saturation magnetization of  $Fe_3O_4$  in  $Fe_3O_4@C$  is 80.4 emu/g.

In this work, to explore the influence of shell modification, the resorcinol/formaldehyde resin (RF) and SiO<sub>2</sub> modified magnetic particles



**Fig. 1.** Synthesis procedure sketch of the  $Fe_3O_4@C$  (a), transmission electronic microscopy (TEM) and scanning electron microscope (SEM) images of  $Fe_3O_4$  (b, c),  $Fe_3O_4@PDA$  (d, e),  $Fe_3O_4@C$  (f, g). (inset) Corresponding TEM images with high magnification (the plotting scale is 60 nm).



Fig. 2. X-ray photoelectron spectroscopy (XPS) of  $Fe_3O_4$ ,  $Fe_3O_4$ @PDA and  $Fe_3O_4$ @C (a), C 1s core-level spectra of  $Fe_3O_4$ @PDA and  $Fe_3O_4$ @C (b), N 1s core-level spectra of  $Fe_3O_4$ @PDA and  $Fe_3O_4$ @C (c).

as contrast samples are also synthesized. A series of characterizations of  $Fe_3O_4@RF$  and  $Fe_3O_4@SiO_2$  particles are carried out, including TEM, SEM, XPS, TG and magnetization hysteresis loops (Figs. S3–S5). Meanwhile, the mechanical properties of  $Fe_3O_4@RF$  and  $Fe_3O_4@SiO_2$ -based magnetic fluids are investigated and the relevant descriptions are in the supplementary material (Fig. S6).

## 2.2. Magneto-responsive rheological characteristics of magnetic fluids based on different magnetic core-shell nanospheres

Magnetic fluids are prepared by dispersing  $Fe_3O_4$ ,  $Fe_3O_4$ @PDA and  $Fe_3O_4$ @C nanospheres into deionized water, respectively. Here, the weight ratio of the magnetic particles is all kept at 10 wt%. The rheological activities of magnetic fluids are investigated by commercially available rheometer (Physia MCR302, Anton Paar) and the magnetic field is applied *via* the magnetorheological unit. Fig. 4a–c show the



Fig. 3. Raman spectra (a), Thermogravimetry (TG) curves (b), X-Ray Diffraction (XRD) patterns (c) and Magnetization hysteresis loops (d) of Fe<sub>3</sub>O<sub>4</sub>, Fe<sub>3</sub>O<sub>4</sub>@PDA and Fe<sub>3</sub>O<sub>4</sub>@C.

typical shear stress vs shear rate curves of as-prepared magnetic fluids. When the applied magnetic field is 0 mT, the shear stress of as-prepared magnetic fluid increases upon increasing the shear rate, which is consistent with the Newtonian fluid (Fig. 4a-c). Under applying constant magnetic field, the shear stress increases with increasing of the shear rate which is similar to traditional magnetic fluids (Fig. 4a-c). Additionally, under applying the external magnetic field, the shear stress of magnetic fluids is substantially enhanced, which is ascribed to the formation of chain-like structure aligned in the magnetic direction. As shown in Fig. 4a and b, the shear stress of Fe<sub>3</sub>O<sub>4</sub> and Fe<sub>3</sub>O<sub>4</sub>@PDA-based magnetic fluid increases with increasing of magnetic flux density from 0 to 110 mT. When the magnetic flux density reaches to 220 mT, the shear stress of these two magnetic fluids are reduced. But the shear stress of Fe<sub>3</sub>O<sub>4</sub>@PDA-based magnetic fluid declines slowly. For Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid, the shear stress values decrease until the magnetic field attains to 440 mT (Fig. 4c). The decline of shear stress values is possibly due to the agglomeration of nanospheres, indicating that the Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid is more stable.

Furthermore, the larger magnetic field usually results in a higher shear stress. Fig. 4d shows the magnetic field dependent shear stress of magnetic fluids composed of Fe<sub>3</sub>O<sub>4</sub>, Fe<sub>3</sub>O<sub>4</sub>@PDA and Fe<sub>3</sub>O<sub>4</sub>@C nanospheres, respectively. Clearly, the shear stress rapidly increases with the enhancement of magnetic field at initial phase, proving the characteristic MR effects. It is apparent that the Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid possesses higher shear stress than other magnetic nanospheres. The Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid is nearly saturated at 180 mT while Fe<sub>3</sub>O<sub>4</sub> and Fe<sub>3</sub>O<sub>4</sub>@PDA based magnetic fluids are almost saturated at 60 mT. The magnetic induced shear stress values for Fe<sub>3</sub>O<sub>4</sub>, Fe<sub>3</sub>O<sub>4</sub>@ PDA and Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid are calculated to be 6.6, 6.4 and 24.9 Pa, respectively (Fig. 4d). In addition, the shear stress decreases with further increasing the magnetic field, which is in accordance with the information exhibited in Fig. 4a-c. The interaction force between magnetic dipoles increase with the increasing of the magnetic field. While the magnetization intensity of particles is saturated, the interaction force will no more increment with the magnetic field increased. Therefore, the shear stress increases quickly when the magnetic field is lower than saturation magnetization. Besides, the mass fraction of all samples is kept at 10 wt% which is quite low, and the particles might tend to assemble in the center of the plate with the application of external magnetic field. Furthermore, because of the low viscosity of water, it is inevitable to exist sealing effect which lead to a reduction in the number of the magnetic field-induced particles chain at high magnetic field, hence the shear stress is decreased.

In previous works, various particles were used for preparing magnetic fluids. In comparison to Fe<sub>3</sub>O<sub>4</sub> nanospheres, the Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid shows preferable MR effect. It is well known that MR effect increases with increasing the particle size. The magnetic induced shear stress value of Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid (24.9 Pa) is even higher than that of MRF-200 (12 Pa), in which the magnetic fluid consists of 200 nm Fe<sub>3</sub>O<sub>4</sub> nanospheres [29]. In comparison to hollow magnetite ferrofluids, the Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid also possesses superior MR property, which is higher than the Fe<sub>3</sub>O<sub>4</sub> hollow spheres and chain-based magnetic fluid [26]. However, when the diameter of hollow spheres rises to 300 nm, the MR effect of hollow nanospheres is stronger [28]. To sum up, the decoration of carbon layer endows magnetic particles with stronger MR performance. Here, the Fe<sub>3</sub>O<sub>4</sub>@C particles also equip with abundant porous structure and high surface area, which help the magnetic fluid possesses high stability. Furthermore, the carbon shell can serve as protective layer, which is beneficial to restrain the oxidization of magnetic particles.

Particle level dynamic simulations are carried out to investigate the correlation between microstructures and MR effect of magnetic fluid. Typical rearrangement of Fe<sub>3</sub>O<sub>4</sub>@C particles in steady shear flow is firstly discussed using the OVITO software [36] (Fig. 5a). In this section, a steady uniform magnetic field of B = 200 mT is chosen as an example. Such magnetic field can ensure the magnetic fluid exhibit the maximum MR effect. The shear rate and weight fraction of magnetic fluid are set at 10 s<sup>-1</sup> and 10 wt%, corresponding to the rheological



Fig. 4. Shear stress vs shear rate curves of magnetic fluids based on  $Fe_3O_4$  (a),  $Fe_3O_4@PDA$  (b) and  $Fe_3O_4@C$  (c) nanospheres, respectively. Magnetic field dependent shear stress of magnetic fluids composed of corresponding  $Fe_3O_4$ -based nanospheres (d).

tests. At the initial state, Fe<sub>3</sub>O<sub>4</sub>@C particles are randomly distributed in the simulation box. After applying the external magnetic field and shear flow for 0.2 ms, particles rapidly form short chains along the field direction due to the magnetic dipolar forces. At 0.4 ms, short chains arrange head-to-tail to form longer chain-like microstructures. With the development of shear flow, chain-like microstructures then incline toward different directions at 10 ms. Dip angles of chains are quite small. In such configuration, the x-component of magnetic dipolar forces are balanced with the hydrodynamic force from the matrix. When the magnetic fluid reaches dynamic equilibrium (40 ms), chain-like microstructures slightly move around the original position due to the Brownian force. Those chains don't flow with the matrix or further aggregate into other particular microstructures during the simulation. Microstructures in dynamic equilibrium under different external fields are plotted in Fig. 5b. Chain-like microstructures possess a consistent slope toward the shear direction (x-axis) under a weak magnetic field



Fig. 5. Evolution of microstructures of  $Fe_3O_4@C$  particles under a magnetic field of 200 mT (a), Microstructures of  $Fe_3O_4@C$  particles at the end of the simulations under different magnetic fields (b).



**Fig. 6.** Average chain Scale and average dip angle of microstructures of  $Fe_3O_4@C$  magnetic fluid at 0.1, 0.2, 0.4, 10, and 40 ms under a 200 mT magnetic field (a); at 40 ms under a magnetic of 20, 60, 100, 140, and 180 mT (b). The lines are plotted to guide the eyes. The weight fraction of the particle is 10 wt%.

(20 mT). The dip angle of chains decreases and become irregular with increasing the external field. If the magnetic field strength exceeds 140 mT, microstructures show familiar characteristics. The x-component of magnetic dipolar forces is in positive correlation with both the magnetic flux density and the inclination of microstructures. Under the same shear rate, the stronger external magnetic field, the small dip angle of chains is required to resist the shear flow.

In the simulations, if the distance between two particles was smaller than 1.2*d* (d = 100 nm), these two particles are considered in the same chain. The dip angle of the chain is defined as the angle between the axis of chain and the external field. The average scale of chains increases from 2.3 to 32.9 particles with the development of shear flow, while the average dip angle increases from 0.99 to 2.51 degrees under a 200 mT magnetic field (Fig. 6a). At the initial state, most of the particles are isolated. The dip angle of chains is meaningless at this time. After applying the magnetic field and shear flow, isolate particles and large chain-like microstructures always exist at the same time. The microstructures are usually not single-particle-width chains. The disparity between the biggest and smallest chains is quite large. The largest chain contains 79 particles and the dip angle varies from 0.23 to 12.53 degrees at the end of the simulation. The average scale of chains at the end of the simulation increases and then declines when increasing the external field (Fig. 6b). The maximum average scale was 46 particles. The average scale under 20 and 180 mT were approximately the same (Fig. 6b). The dip angle of chains exhibits a decreasing trend from 6.0 to 3.4 degrees with increasing the magnetic field, which is consistent with the snapshot of microstructures (Fig. 5).

Meanwhile, the simulated shear stress as a function of magnetic flux density is also discussed (Fig. 7). Shear stress first linearly increases with the external field and reaches saturation after B = 150 mT. A good correspondence between simulations and experiments is observed, which confirms the particle-level simulation method has a high accuracy. Simulated shear stress of magnetic fluid based on Fe<sub>3</sub>O<sub>4</sub>@C particles is also stronger than the experimental shear stress of pure Fe<sub>3</sub>O<sub>4</sub> magnetic fluid. The microstructure evolution of Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid is similar the one prepared by Fe<sub>3</sub>O<sub>4</sub> particles. The improvement of magnetic moment of Fe<sub>3</sub>O<sub>4</sub>@C particles induces larger shear stress than the Fe<sub>3</sub>O<sub>4</sub> particles.

The changing processes of magnetic fluids based on  $Fe_3O_4$  and  $Fe_3O_4@C$  nanospheres from a random status to chain-like structure are displayed in Fig. S7. When an external magnetic field is applied, the random dispersed particles in magnetic fluid assemble to form chain-like structure along the direction of magnetic field. With the exerting of steady shear, the chains tilt with quite small dip angles. This process is depended on the acting force between the magnetic particles. Fig. 8a presents the force analyses of magnetic fluids based on  $Fe_3O_4$  and  $Fe_3O_4@C$  nanospheres, respectively. The  $Fe_3O_4@C$  possesses higher



Fig. 7. Shear stress as a function of magnetic flux density, a comparison between experiments and simulations.

crystallinity and shows bigger magnetic torque, which leads to stronger magnetic dipolar forces. The enhanced magnetic dipolar forces and unique properties of Fe<sub>3</sub>O<sub>4</sub>@C greatly strengthen the as-formed chain-like structures, thus the Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid has a larger MR effect than that of Fe<sub>3</sub>O<sub>4</sub>.

To further verify the inference, the actual chain-like structure of magnetic fluid are observed in a Poiseuille flow. As shown in Fig. 8c, the microfluidic channel is manufactured in a slice of poly-dimethylsiloxane (PDMS) with a 280  $\mu$ m × 120  $\mu$ m rectangular section. A permanent magnet is parallelly placed to the channel to generate a homogeneous magnetic field. The volume flow is set to 0.5 ml/h and the frame rate of high-speed camera is 20,000 fps. From Fig. 8b it can be seen that Fe<sub>3</sub>O<sub>4</sub> and Fe<sub>3</sub>O<sub>4</sub>@C-based magnetic fluid both form chains-like structure parallel to the magnetic field, which are consistent with the simulation results. In contrast to the Fe<sub>3</sub>O<sub>4</sub>, the Fe<sub>3</sub>O<sub>4</sub>@C fluid (Fig. 8b (II)) shows thicker and sparse particle chains structure, which is attributed to stronger interaction forces between Fe<sub>3</sub>O<sub>4</sub>@C particles. So the Fe<sub>3</sub>O<sub>4</sub>@C based magnetic fluid exhibits much better MR properties.

#### 3. Conclusion

In this work, a novel kind of  $Fe_3O_4@C$  magnetic particles are prepared by surface modification. These  $Fe_3O_4@C$  particles are dispersed into water to prepare magnetic fluid and its rheological properties are measured.  $Fe_3O_4@C$ -based magnetic fluid shows better mechanical properties than the pristine  $Fe_3O_4$  and  $Fe_3O_4@PDA$  particles. It is found



Fig. 8. Force analyses of magnetic fluid based on  $Fe_3O_4$  and  $Fe_3O_4$ @C nanospheres (a), Snapshots of magnetic fluid based on  $Fe_3O_4$  (I) and  $Fe_3O_4$ @C (II) in Poiseuille flow in a microfluidic channel (b), schematic diagram of the microfluidic channel and experimental setup (c).

that the carbon layer and unique porous inner-structure of  $Fe_3O_4@C$  play critical roles in determining the MR effects of the final magnetic fluid. Furthermore, the Particle level dynamic simulations about microstructure evolution and magnetic field dependent shear stress are carried out. The chain-like model well explains the MR mechanism and the simulation results matches perfectly with the experimental data. Due to the stronger magnetic dipolar forces, chain-chain interaction and much steadier chain-like structure, the  $Fe_3O_4@C$ -based magnetic fluid exhibits enhanced MR effects. This work is valuable for researcher to further prepare high performance magnetic fluid and understand the origination of MR effects.

#### CRediT authorship contribution statement

Linfeng Bai: Investigation, Formal analysis, Resources, Writing original draft. Lei Pei: Software. Saisai Cao: Investigation. Xiaokang He: Investigation. Min Sang: Writing - review & editing. Wanquan Jiang: . Shouhu Xuan: Conceptualization, Writing - review & editing, Supervision. Xinglong Gong: .

#### **Declaration of Competing Interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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#### Appendix A. Supplementary data

Supplementary data to this article can be found online at https://doi.org/10.1016/j.jmmm.2020.166734.

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